

1 **Long-Term Rainfall Risk from Tropical Cyclones in Coastal Areas**

2 By

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32 **Abstract**

33 We develop a methodology for the frequency of extreme rainfall intensities caused by tropical
34 cyclones (TCs) in coastal areas. The model does not account for landfall effects. This makes the
35 developed framework best suited for open-water sites and coastal areas with flat topography. The
36 mean rainfall field associated with a TC with maximum tangential wind speed V_{max} , radius of
37 maximum winds R_{max} , and translation speed V_t is obtained using a physically-based model
38 (Langousis and Veneziano, 2008), whereas rainfall variability at both large scales (from storm to
39 storm) and small scales (due to rainbands and local convection) is modeled statistically. The
40 statistical component is estimated using precipitation radar (PR) data from the TRMM mission.
41 Taylor's hypothesis is used to convert spatial rainfall intensity fluctuations to temporal
42 fluctuations at a given location A. The combined physical-statistical model gives the distribution
43 of the maximum rainfall intensity at A during an averaging period D for a TC with
44 characteristics (V_{max}, R_{max}, V_t) that passes at a given distance from A. To illustrate the use of the
45 model for long-term rainfall risk analysis, we formulate a recurrence model for tropical cyclones
46 in the Gulf of Mexico that make landfall between longitudes 85° - 95° W. We then use the rainfall
47 and recurrence models to assess the rainfall risk for New Orleans. For return periods of 100 years
48 or more and long averaging durations (D around 12-24 hours), tropical cyclones dominate over
49 other rainfall event types, whereas the reverse is true for shorter return periods or shorter
50 averaging durations.

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52 **Keywords:** Rainfall Extremes, IDF Curves, Tropical Cyclones, Tropical Meteorology, Floods.

1. Introduction

The quantification of long-term rainfall risk is a basic problem of stochastic hydrology (e.g. Chow *et al.*, 1988; Singh, 1992). Our specific interest is in the risk of extreme rainfall posed at coastal sites by tropical cyclones (TCs). These events are relatively rare, but in combination with wind, surge and waves, high rainfall intensities may have devastating consequences (Herbert *et al.*, 1997; Rappaport, 2000).

For ordinary rainfall, standard risk analysis techniques use historical annual-maximum data (e.g. Koutsoyiannis *et al.*, 1998) or peak-over-threshold (PoT) information (e.g. Madsen *et al.*, 1997). The episodic and spatially localized nature of tropical cyclones prevents one from using these standard techniques. For example, the annual maximum and PoT rainfall statistics due to tropical cyclones are very sensitive to whether the site is “hit” by one or more TCs during a year and therefore are highly erratic. For this reason, the risk is best assessed parametrically, by combining a probabilistic model of the maximum rainfall due to a TC with given characteristics $\theta = [\theta_1, \dots, \theta_r]$ with the rate at which those events occur. For coastal sites, the vector θ might include the intensity and size of the storm, the location and translational velocity at landfall, and possibly other parameters related to atmospheric conditions, the radial profile of the tangential winds, etc. Parametric approaches of this type have been used to assess the risk posed by tropical cyclones for wind, surge and waves (Myers, 1975; Ho and Myers, 1975; Ho *et al.*, 1987; Powell *et al.*, 2005; IPET, 2006, 2008), but not rain. Here we develop a parametric approach to calculate peak rainfall intensities from tropical cyclones, and use this approach to study the importance of TCs relative to other storm types and determine the TC characteristics that dominate different levels of risk.

77 The main problem for rainfall is to evaluate the extreme precipitation intensities caused by a
78 TC with given characteristics θ . The historical data are too sparse and the potentially important
79 TC parameters are too many to infer such extreme rainfalls from empirical observations alone.
80 For example, current empirical approaches (Lonfat *et al.*, 2004, 2007; Tuleya *et al.*, 2007)
81 classify storms into three coarse intensity categories and use microwave imager (TMI) data from
82 TRMM (Simpson *et al.* 1988) to calculate the ensemble-average rainrate for each category as a
83 function of distance from the TC center.

84 The alternative we pursue here is to use a physical model to assess the dependence of the
85 mean rainfall field on θ and statistical analysis to quantify the fluctuations of rainfall intensity
86 around this mean field. The physical model is that developed by Langousis *et al.* (2008) and
87 Langousis and Veneziano (2008). Langousis *et al.* (2008) proposed a theoretical method to
88 estimate the large-scale horizontal and vertical winds inside TCs (the vertical winds are largely
89 responsible for rain). The model is an extension of Smith's (1968) formulation and is referred to
90 here as the Modified Smith (MS) model. Characteristics of the TC that are explicitly considered
91 by the model are the maximum tangential wind speed V_{max} , the radius of maximum winds R_{max} ,
92 the parameter B that controls the shape of the radial profile of the tangential wind speed
93 (Holland, 1980), the storm translation velocity V_t , the surface drag coefficient C_d , and the vertical
94 diffusion coefficient K . When $V_t = 0$, the wind field is symmetric around the storm center,
95 whereas when the TC translates in the Northern (Southern) hemisphere the field is asymmetric,
96 with stronger horizontal and vertical winds right-front (left-front) of the storm. The model does
97 not resolve rainbands, local convection and turbulent phenomena and therefore produces smooth
98 wind fields.

99 Langousis and Veneziano (2008) extended the MS model to predict TC rain, assuming that
100 the upward moisture flux at the top of the TC boundary layer is all converted into rainfall. The
101 vertical moisture flux is evaluated from the vertical winds generated by the MS model and two
102 additional parameters: the average temperature \bar{T} and average saturation ratio \bar{Q} inside the TC
103 boundary layer. We call this the modified-Smith-for-rainfall (MSR) model. The MSR model
104 should prove useful for climatologic studies, but for hazard analysis it has the limitation of
105 ignoring the inter-storm and intra-storm variations of rainfall intensity. These variations are
106 highly significant for the assessment of risk. For example, Lonfat *et al.* (2004) found that, also
107 within a given TC strength category, the average of the positive rainfall intensity inside annular
108 regions of 10 km width may deviate from the median value by more than one order of
109 magnitude.

110 Our main objectives are: (1) Extend the MSR model to obtain the probability distribution of
111 the maximum rainfall intensity in an averaging time interval of given duration D at a fixed
112 geographical location during the passage of a tropical cyclone with given characteristics θ , and
113 (2) Combine this maximum rainfall model with a TC recurrence model to quantify rainfall risk in
114 the form of intensity-duration-frequency (IDF) curves. For the first objective, we consider a site
115 A at some distance y to the right ($y < 0$) or left ($y > 0$) of the moving TC center, as shown in
116 Figure 1. As the storm passes, the rainfall intensity at A fluctuates as a random process $I(t)$. Our
117 interest is in $I_D(t)$, the moving average of $I(t)$ for an averaging duration D , and more specifically
118 in the distribution of $I_{D,max}(y,\theta)$, the maximum of $I_D(t)$ during the storm.

119 Section 2 presents our general approach to calculate the distribution of $I_{D,max}(y,\theta)$. This
120 distribution is obtained in Section 3 and validated in Section 4. Section 4 also shows how the
121 distribution depends on various storm characteristics, the standardized distance y/R_{max} from the

122 center of the storm, and the averaging duration D . Section 5 uses the model of $I_{D,max}(y,\theta)$ and a
123 recurrence relation for hurricanes in the Gulf of Mexico to obtain IDF curves for New Orleans
124 and compares these curves with published IDF values for all rainstorms (TCs and non-TCs)
125 combined. Conclusions are stated in Section 6.

126 **2. A Framework for the Estimation of Extreme TC Rainfall**

127 Our first objective is to relate the distribution of the maximum rainfall intensity $I_{D,max}(y,\theta)$ to the
128 smooth rainfall intensities produced by the MSR model of Langousis and Veneziano (2008). The
129 storm parameters are $\theta = [V_{max}, R_{max}, V_t]$. The analysis uses a Cartesian reference frame (x,y) ,
130 translated and rotated such that the center of the storm O moves to the right along the x axis; see
131 Figure 1. In this reference, the ordinate y of A is also the closest (signed) distance of A from the
132 storm center.

133 To estimate this relationship, we use precipitation radar (PR) data from the TRMM mission
134 (Simpson *et al.* 1988; Kummerow *et al.*, 1998; Lee *et al.*, 2002). These data are in the form of
135 swaths about 200km wide with a spatial resolution of approximately 5 km and have been
136 validated against ground-based radar and rain gauge measurements (Bolen and Chandrasekar,
137 2000; Liao *et al.*, 2001; Wolff *et al.*, 2005). Due to their long inter-frame time (about 12 hours),
138 the PR snapshots cannot be interpolated to produce the rainfall intensities in continuous time that
139 are needed to estimate rainfall maxima. A common way to overcome this limitation is to use
140 Taylor's frozen turbulence hypothesis (Taylor, 1921, 1938). Under this hypothesis, the temporal
141 variability of rainfall at a fixed location A is statistically the same as the variability that results
142 from translating the frozen-in-time rainfield over A with the storm velocity V_t . For example,
143 Vicente *et al.* (1998), Scofield and Kuligowski (2003), Kidder *et al.* (2005) and Ferraro *et al.*

144 (2005) used Taylor's hypothesis to obtain rainfall totals at fixed locations from satellite and radar
 145 rainfall snapshots.

146 It follows from Taylor's hypothesis that $I_{D,max}(y, \boldsymbol{\theta})$ has the same distribution as $I_{l,max}(y, \boldsymbol{\theta})$, the
 147 maximum of the rainfall intensity averaged in a spatial window of length l along cross-section C
 148 in Figure 1, for $l = DV_r$. As an example, Figure 2 shows moving-average rainfall intensities from
 149 Hurricane Katrina (2005) along a cross-section at distance $y = 100$ km from the storm center, for
 150 averaging lengths $l = 6$ km (dashed line) and $l = 24$ km (solid line). The cross-section extends
 151 over $L = 384$ km and is symmetrical relative to the storm center.

152 The intensity labeled I_L in Figure 2 is the average PR rainrate in L , whereas $I_{L,MSR}$ is the
 153 estimate of that average rainrate produced by the MSR model. These average intensities play an
 154 important role in our analysis. For any given $(y, \boldsymbol{\theta})$ combination, the model estimate $I_{L,MSR}$ is
 155 fixed, whereas I_L is regarded as a random variable with different values for different tropical
 156 cyclones. We model this storm-to-storm variability by expressing $I_L(y, \boldsymbol{\theta})$ as

$$157 \quad I_L(y, \boldsymbol{\theta}) = I_{L,MSR}(y, \boldsymbol{\theta}) \beta_L(y, \boldsymbol{\theta}) \quad (1)$$

158 where β_L is a random variable.

159 Figure 2 also shows significant amplification of the rainfall intensity when one considers the
 160 maximum over lengths $l < L$. One may express the maximum in l , $I_{l,max}$, as

$$161 \quad I_{l,max}(y, \boldsymbol{\theta}) = I_{L,MSR}(y, \boldsymbol{\theta}) \beta_{l,max}(y, \boldsymbol{\theta}) = I_{L,MSR}(y, \boldsymbol{\theta}) \beta_L(y, \boldsymbol{\theta}) \gamma_{l,max}(y, \boldsymbol{\theta}) \quad (2)$$

162 where the total factor relative to $I_{L,MSR}$, $\beta_{l,max}$, is the product of β_L in equation (1) and a random
 163 amplification factor $\gamma_{l,max}$ for the change of scale from L to l . The next section uses PR/TRMM
 164 data from 8 tropical cyclones (a total of 38 frames) to derive the distributions of β_L and $\gamma_{l,max}$. The
 165 selected frames (see Table 1) cover a wide range of TC intensities, from tropical storms to CAT5
 166 systems, under pre-landfall conditions. This makes our model best suited for open-water sites,

167 but it should also be accurate in coastal areas with a flat topography. For example, Marks *et al.*
 168 (2002) (see also Tuleya *et al.*, 2007) used TMI rainfall products for TCs over water to predict
 169 rainfall rates at inland locations. For sites close to the shore, the predictions had low bias relative
 170 to raingauge measurements.

171 Due to the limited lateral coverage of the PR instrument, an additional requirement for
 172 selecting the frames was to cover regions close to the hurricane core (with radial distance less
 173 than 300 km from the storm center), as these are the regions that are most critical for rainfall.

174 **3. Distribution of β_L and $\gamma_{l,max}$**

175 Equation (2) relates the maximum rainfall intensity in l to the average intensity in L produced by
 176 the MSR model using two random factors: a factor β_L to obtain the average rainfall in L , and a
 177 factor $\gamma_{l,max}$ to obtain the maximum average intensity at a smaller scale l . Sections 3.1 and 3.2
 178 obtain the distribution of these factors using the rainfall information in Table 1 and MSR model
 179 simulations.

180 **3.1 Distribution of β_L**

181 The factor β_L is given by

$$182 \quad \beta_L(y, \boldsymbol{\theta}) = \frac{I_L(y, \boldsymbol{\theta})}{I_{L,MSR}(y, \boldsymbol{\theta})} \quad (3)$$

183 where I_L and $I_{L,MSR}$ are the same as in equation (1). The distribution of β_L generally depends on
 184 the distance y from the TC center and the vector $\boldsymbol{\theta} = [V_{max}, R_{max}, V_l]$ of storm characteristics, but
 185 as we show next, a simple parameterization in terms of the standardized distance $y' = |y/R_{max}|$ and
 186 the large-scale MSR rainfall intensity $I_{L,MSR}$ suffices. Of course, $I_{L,MSR}$ is itself a function of $\boldsymbol{\theta}$.

187 Figure 3 shows statistics of β_L as a function of y' and $I_{L,MSR}$ for the TRMM frames in Table 1.
 188 For each frame, the $I_{L,MSR}$ intensities at different distances y from the center of the storm were

189 calculated using the MSR model and the values of V_{max} , R_{max} , and V_t in the extended best track
 190 record (M. DeMaria, 2008; personal communication; Demuth *et al.*, 2006); see Table 1. In
 191 addition to V_{max} , R_{max} , and V_t , the MSR model requires specification of the vertical diffusion
 192 coefficient K , the surface drag coefficient C_d , the vertically averaged temperature \bar{T} and
 193 saturation ratio \bar{Q} inside the boundary layer (BL), Holland's B parameter for the profile of
 194 gradient winds, the sloping angle ψ_0 and height H_0 of the wall updraft, and the temporal scale t_r
 195 for azimuthal re-distribution of rainfall by the cyclonic circulation; see Langousis and Veneziano
 196 (2008) for details. In our simulations we have set $K = 50 \text{ m}^2/\text{s}$, $C_d = 0.002$, $\bar{T} = 22^\circ\text{C}$, $\bar{Q} = 0.8$,
 197 $B = 1$, $\psi_0 = 50^\circ$, $H_0 = 6 \text{ km}$ and $t_r = 60 \text{ min}$. Langousis and Veneziano (2008) recommend these
 198 settings as representative of tropical cyclones in the North Atlantic and as values that reproduce
 199 well the TRMM/PR rainfall fields in an ensemble-average sense.

200 Figures 3.a and 3.b show smooth contour plots of the log-mean $m_{\ln\beta_L}$ and log-standard
 201 deviation $\sigma_{\ln\beta_L}$ of β_L as a function of the standardized distance $y' = |y/R_{max}|$ and the MSR rainfall
 202 intensity $I_{L,MSR}$ for the 38 frames in Table 1. For each frame, a regular spacing $\Delta y = 10 \text{ km}$ was
 203 maintained between adjacent transects, producing a total of 789 points; see Figure 3.a. In all
 204 cases, averaging is over segments of length $L = 384 \text{ km}$, symmetric relative to the storm center.
 205 This value of L encompasses more than 95% of the total rainfall volume along each transect; see
 206 for example Figure 2. Smooth estimates of the mean value and variance of $\ln\beta_L$ were obtained
 207 using an isotropic Gaussian kernel with standard deviation 0.5 in the $[\ln(I_{L,MSR}), \ln(y')]$ -plane.
 208 Hence, if $g(\mathbf{x})$ denotes this kernel, local estimates of $m_{\ln\beta_L}(\mathbf{x}_0)$ and $\sigma_{\ln\beta_L}^2(\mathbf{x}_0)$ around a given point
 209 $\mathbf{x}_0 = [\ln(I_{L,MSR}), \ln(y')]$ are given by

$$\begin{aligned}
210 \quad m_{\ln\beta_L}(\mathbf{x}_0) &= \frac{\sum_i \ln\beta_L(\mathbf{x}_i) g(\mathbf{x}_i - \mathbf{x}_0)}{\sum_i g(\mathbf{x}_i - \mathbf{x}_0)}, \quad \sigma_{\ln\beta_L}^2(\mathbf{x}_0) = \frac{\sum_i [\ln\beta_L(\mathbf{x}_i) - m_{\ln\beta_L}(\mathbf{x}_0)]^2 g(\mathbf{x}_i - \mathbf{x}_0)}{\sum_i g(\mathbf{x}_i - \mathbf{x}_0)} \quad (4)
\end{aligned}$$

211 where \mathbf{x}_i is the generic $[\ln(I_{L,MSR}), \ln(y')]$ combination for which a value of β_L is available. To use
212 values of β_L at locations close to the center of the storm where $\ln(y')$ diverges, 59 points with
213 $|y| < 0.5 R_{max}$ where moved to $y = 0.5 R_{max}$.

214 The overall mean value of β_L is 1.02, indicating that on average the MSR model produces
215 unbiased large-scale estimates of the PR rainrates. The dashed lines in Figure 3.b delimit the
216 region of high data density and are generally oriented along the gradient of $\sigma_{\ln\beta_L}$. Figure 3.c
217 shows plots of $m_{\ln\beta_L}$ and $\sigma_{\ln\beta_L}$ as a function of the transformed variable $\omega = \ln(y') - 0.4\ln(I_{L,MSR})$
218 along the dashed-dotted line in Figure 3.b. The log-mean $m_{\ln\beta_L}$ is approximately constant and
219 equal to -0.5, whereas $\sigma_{\ln\beta_L}$ increases as the standardized distance y' increases or the large-scale
220 mean rainfall intensity $I_{L,MSR}$ decreases. This higher log variability in regions of lower intensity is
221 expected due to the more episodic nature of rainfall in those regions. This is also in qualitative
222 agreement with the findings of Lonfat *et al.* (2004) and Molinari *et al.* (1994). The solid lines in
223 Figure 3.c are least-squares fits for the mean and standard deviation of $\ln\beta_L$. For y close to zero,
224 the fitted standard deviation becomes very small or negative. To avoid this inconsistency, we
225 have imposed a lower bound of 0.5 to the fitted standard deviation.

226 To investigate the distribution type, we standardize the empirical values of $\ln\beta_L$ by removing
227 the parametrically fitted mean -0.5 and dividing by the parametrically fitted standard deviation
228 $0.25\omega + 0.87$. Figure 3.d shows a histogram of these standardized quantities and suggests that
229 $\ln\beta_L$ has near-normal distribution. To check for possible lack of fit and possible dependence of

230 $\ln\beta_L$ on other parameters, we generated histograms of the type in Figure 3.d separately for
 231 different ranges of y , $I_{L,MSR}$, R_{max} and V_{max} ; see Langousis (2008). As none of these analyses
 232 reveals significant dependence, we use the fits in Figure 3.c and model $\ln\beta_L$ as a normal variable
 233 with parameters

$$\begin{aligned}
 234 \quad m_{\ln\beta_L}(\omega) &= -0.5 \\
 \sigma_{\ln\beta_L}(\omega) &= \max\{0.5, 0.25\omega + 0.87\}
 \end{aligned} \tag{5}$$

235 where $\omega = \ln(y') - 0.4\ln(I_{L,MSR})$.

236 **3.2 Distribution of $\gamma_{l,max}$**

237 Next we consider the amplification factor $\gamma_{l,max}$ in equation (2). The distribution of this factor can
 238 be found by a variety of methods, from the direct use of data on $\gamma_{l,max}$ from the frames in Table 1
 239 to theoretical analysis of the maximum of the moving-average processes $I_l(x)$ illustrated in Figure
 240 2. Langousis (2008) compared several such approaches and found similar results. Here we
 241 follow the empirical approach, which is the simpler and more transparent method. We start by
 242 calculating the empirical ratio

$$243 \quad \gamma_{l,max} = \frac{I_{l,max}}{I_L}, \quad l \leq L \tag{6}$$

244 where I_L is the average PR rainrate along a cross section C of fixed length $L = 384$ km and $I_{l,max}$
 245 is the maximum rainfall intensity when the same cross section is continuously scanned using an
 246 averaging window of length l ; see Figures 1 and 2 and Section 2. Ideally, the cross section C
 247 should be in the direction of the storm motion, but since the TRMM swaths are not always
 248 aligned with that direction, we calculate the factor $\gamma_{l,max}$ using cross-sections parallel to the swath
 249 track. Hence, the resulting factor $\gamma_{l,max}$ does not depend on the orientation of C relative to the
 250 storm motion. Langousis (2008) verified that $\gamma_{l,max}$ is insensitive to this orientation by dividing

251 the swaths into two groups: those that are generally aligned with the storm trajectory and those
 252 that are not. The distribution of $\gamma_{l,max}$ is similar in the two cases.

253 Langousis (2008) also studied the dependence of the distribution of $\gamma_{l,max}$ on R_{max} . Dependence
 254 is expected because smaller values of R_{max} produce more peaked radial rainfall profiles and hence
 255 higher rainfall maxima. The finding is that for small spatial scales ($l \leq 12$ km) the mean value
 256 and standard deviation of $\gamma_{l,max}$ increase somewhat with decreasing R_{max} , whereas at larger spatial
 257 scales the increase is modest. Based on these results, we ignore the dependence of $\gamma_{l,max}$ on R_{max}
 258 and use a simple parameterization in terms of the averaging length l and the large-scale average
 259 intensity I_L . The latter quantity depends significantly on both the storm intensity V_{max} and the
 260 distance y from the storm center; see Langousis and Veneziano (2008).

261 Figure 4 shows log-log plots of $E[\gamma_{l,max}]$ and $\text{Var}[\gamma_{l,max}]$ against l after classifying the 789
 262 cross-sections in Figure 3.a into 12 equally-sized I_L bins. As expected, $\text{Var}[\gamma_{l,max}]$ increases with
 263 decreasing spatial scale l . A less obvious finding is that the variability of $\gamma_{l,max}$ increases as the
 264 large-scale intensity I_L decreases. Considering that lower values of I_L are generally found at
 265 larger distances y from the storm center, Figure 4 shows that the outer TC environment exhibits
 266 higher (multiplicative) variability relative to the inner region. The higher variability inside low- I_L
 267 regions is due for the most part to an increase in the dry area fraction (Langousis, 2008) and has
 268 been noted also in other studies (Molinari *et al.*, 1994; Lonfat *et al.*, 2004). This feature is also
 269 commonly observed in extra-tropical rainfall (e.g. Over and Gupta, 1996; Deidda *et al.*, 2006;
 270 Veneziano *et al.*, 2006a; Gebremichael *et al.*, 2006).

271 For each intensity category I_L , we use least squares to fit linear and quadratic expressions for
 272 the log-mean and log-variance of $\gamma_{l,max}$,

$$\begin{aligned}
 \ln E[\gamma_{l,max}] &= a_1 \ln l + a_2 \\
 \ln \text{Var}[\gamma_{l,max}] &= a_3 (\ln l)^2 + a_4 \ln l + a_5
 \end{aligned}
 \tag{7}$$

274 where $l \leq L$ is in km and a_1 - a_5 are parameters. Figure 5 shows how the parameters a_1 - a_5 in
 275 equation (7) vary with the large-scale rainfall intensity I_L . The solid lines in Figure 5 are smooth
 276 least-squares estimates of a_i ($i=1,\dots,5$). Use of the smooth estimates reproduces well the
 277 empirical moments of $\gamma_{l,max}$; see solid lines in Figure 4.

278 The amplification factor $\gamma_{l,max}$ has values between 1 and L/l . The lower bound corresponds to a
 279 uniform distribution of rainfall inside L , whereas the upper bound is attained when all the rainfall
 280 in L is concentrated in a single l interval. We model $\gamma_{l,max}$ using a beta distribution with moments
 281 in equation (7). One may write this cumulative distribution as

$$282 \quad F_{\gamma_{l,max}}(\gamma) = F_X\left(\frac{\gamma-1}{L/l-1}\right), \quad \gamma \geq 1 \quad (8)$$

283 where F_X is the beta distribution in $[0,1]$ with parameters

$$284 \quad E[X] = \frac{E[\gamma_{l,max}]-1}{L/l-1}, \quad \text{Var}[X] = \frac{\text{Var}[\gamma_{l,max}]}{(L/l-1)^2} \quad (9)$$

285 Figure 6 compares the empirical distribution of $\gamma_{l,max}$ at spatial scales $l = 96$ and 6 km for
 286 different large-scale average intensities I_L with theoretical distributions from equations (8) and
 287 (9). The moments $E[\gamma_{l,max}]$ and $\text{Var}[\gamma_{l,max}]$ in equation (9) are calculated using equation (7) with
 288 parameters a_1 - a_5 in Figure 5. Equally good fits are obtained for other window sizes l ; see
 289 Langousis (2008).

290 **4. Validation of Maximum Rainfall Model and Sensitivity Analysis**

291 For a tropical cyclone with parameters $\theta = [V_{max}, R_{max}, V_i]$ and a given distance y from the storm
 292 center, one may use equation (2) and the distributions of β_L and $\gamma_{l,max}$ in Section 3 to obtain the
 293 distribution of the maximum rainfall intensity $I_{l,max}$ as

294
$$P[I_{l,max}(y, \boldsymbol{\theta}) \leq i] = \int_0^{\infty} f_{I_L|y, \boldsymbol{\theta}}(u) F_{\gamma_{l,max}|I_L=u}(i/u) du \quad (10)$$

295 where $f_{I_L|y, \boldsymbol{\theta}}$ is the probability density function of $I_L = I_{L,MSR} \beta_L$ given $(y, \boldsymbol{\theta})$ and $F_{\gamma_{l,max}|I_L}$ is the
 296 cumulative distribution function of $\gamma_{l,max}$ given I_L . To assess the validity of the probabilities
 297 generated by equation (10), we compare them with observed relative frequencies, as follows. For
 298 each of the 789 transects extracted from the PR data in Table 1,

- 299 1. We calculate the maximum intensity $I_{l,max}$ over segments of different length l ;
- 300 2. We use (V_{max}, R_{max}, V_t) from Table 1 and the distance y of the transect from the TC center
 301 to obtain model estimates of the large-scale mean rainfall intensity $I_{L,MSR}(y, \boldsymbol{\theta})$ for $L=$
 302 384 km. All other MSR model parameters are fixed to the values in Section 3.1.
- 303 3. We use equation (10) and the parametric expressions in equations (5) and (7) and Figure
 304 5 to find the distribution of $I_{l,max}$ and the probability P with which the observed value
 305 from step (1) is not exceeded.

306 If the model is correct, the probabilities P from step (3) have uniform distribution between 0 and
 307 1. Figure 7 shows histograms of P for different l . One sees that the histograms differ somewhat
 308 from a uniform density (the chi-square goodness of fit test applied to the bins shown in Figure 7
 309 passes at a level of significance around 0.005-0.01 depending on the scale of averaging l). We
 310 have investigated this issue in some detail (Langousis, 2008) and found that the biases are due
 311 mainly to dependence of the amplification factor $\gamma_{l,max}$ on the radius of maximum winds R_{max} ; see
 312 Section 3.2. Although a parameterization of $\gamma_{l,max}$, that includes R_{max} as an independent variable
 313 would improve the goodness of fit, here we retain the simpler model.

314 The distribution of $I_{l,max}$ in equation (10) depends critically on the amplification factor $\beta_{l,max}$ in
 315 equation (2). Figure 8 shows how the distribution of $\beta_{l,max}$ depends on l , V_{max} , and $y' = |y/R_{max}|$.

316 The effect of the translation velocity V_t is modest and is not displayed. Also, for given V_{max} and
 317 $y' = |y/R_{max}|$, $\beta_{l,max}$ is insensitive to R_{max} . The dispersion of $\beta_{l,max}$ increases as l decreases. It also
 318 increases for smaller V_{max} and larger y' . The latter effects are related to the increased spatial
 319 variability of the rainfall intensity in regions of lower average precipitation.

320 **5. Long-term Rainfall Risk for New Orleans**

321 To assess rainfall risk at a given location A, one must find the rate $\lambda_{I_{D,max} > i}$ of tropical cyclones
 322 for which $I_{D,max}$, the maximum rainfall intensity at A for a given averaging duration D , exceeds
 323 different threshold levels i . This rate is given by

$$324 \quad \lambda_{I_{D,max} > i} = \lambda \int_{\text{all } (y, \boldsymbol{\theta})} P[I_{D,max}(y, \boldsymbol{\theta}) > i] f_{y, \boldsymbol{\theta}}(y, \boldsymbol{\theta}) dy d\boldsymbol{\theta} \quad (11)$$

325 where λ is the rate of TCs in the region, $P[I_{D,max}(y, \boldsymbol{\theta}) > i]$ is the probability that, for a storm with
 326 characteristics $\boldsymbol{\theta}$, $I_{D,max}$ at distance y from the storm center exceeds i , and $f_{y, \boldsymbol{\theta}}$ is the joint density
 327 of $(y, \boldsymbol{\theta})$. The joint density $f_{y, \boldsymbol{\theta}}$ and the rate λ are region-specific and define the TC recurrence
 328 model. Under Taylor's hypothesis, $P[I_{D,max}(y, \boldsymbol{\theta}) > i]$ is obtained by setting $l = DV_t$ in equation
 329 (10).

330 To exemplify, we use equation (11) and a recurrence model for an appropriate coastal region
 331 of the Gulf of Mexico to obtain intensity-duration-frequency (IDF) relationships for New
 332 Orleans. We select this location because: 1) the site is close to the coast and has a flat
 333 topography; hence our pre-landfall model should produce accurate results, 2) a number of studies
 334 have developed TC recurrence models for the Louisiana coast, and 3) one can compare the TC
 335 rainfall results with available IDF curves from continuous rainfall records in the region.

336 **5.1 TC recurrence model for the northern Gulf of Mexico**

337 We start by specifying the distribution of the distance y between the center of the storm and the
 338 city of New Orleans (point A), which is located at approximately (90°W, 30°N). Then we
 339 consider the distribution of $\theta = [V_{max}, R_{max}, V_i]$. The joint model for V_{max} and R_{max} is specified
 340 through the distribution of the maximum pressure deficit ΔP_{max} and the conditional distributions
 341 of $[V_{max}|\Delta P_{max}]$ and $[R_{max}|\Delta P_{max}]$. Finally we specify the TC rate λ . To keep the model simple, we
 342 approximate the coastline by a line segment with constant latitude 30°N and longitudinal range
 343 85°-95°W (≈ 960 km), centered at A.

344 Let z be the location (positive eastward) of landfall relative to A. Assuming a straight storm
 345 path, the closest distance of the storm center from the site is

$$346 \quad y = -z \cos(\alpha) \quad (12)$$

347 where α is the azimuth of the storm track at landfall, positive clockwise. The distribution of y
 348 can be obtained numerically from equation (12) and the distributions of α and z , assumed here to
 349 be independent. For z we use a uniform distribution in the interval [85°W, 95°W]. The
 350 distribution of the angle α in the region is usually found to be normal or the mixture of two
 351 normal distributions, one for easterly storms and the other for westerly storms (Vickery and
 352 Twisdale, 1995; IPET, 2006, 2008). Here we model α using a single normal distribution with
 353 mean value $m_\alpha = -5.4^\circ$ and standard deviation $\sigma_\alpha = 34.9^\circ$. This distribution was obtained by IPET
 354 (2006) using NOAA's HURDAT data set (Jarvinen *et al.*, 1984) and found to describe well
 355 storms with central pressure deficit $\Delta P_{max} > 34$ hPa that make landfall in the longitudinal range
 356 85°-95°W.

357 Several studies (Holland, 1980; Atkinson and Holiday, 1977; Willoughby and Rahn, 2004)
 358 have used theoretical arguments and pressure-wind observations to relate V_{max} to ΔP_{max} . The
 359 relationships are typically of the power-law type

360
$$V_{max} = c (\Delta P_{max})^g \quad (13)$$

361 where c and g are positive constants. Using flight level data from 23 hurricane seasons,
 362 Willoughby and Rahn (2004) found $c = 4.8$ and $g = 0.559$ for V_{max} in m/s and ΔP_{max} in hPa.
 363 Based on these and other findings of Willoughby and Rahn (2004), we model $[V_{max} | \Delta P_{max}]$ as a
 364 lognormal variable with mean value $4.8(\Delta P_{max})^{0.559}$ and coefficient of variation 0.15.

365 Empirical evidence (Vickery and Twisdale, 1995; Vickery *et al.*, 2000; Willoughby and Rahn,
 366 2004; Powell *et al.*, 2005; IPET, 2008) and theoretical arguments (Shen, 2006) show that R_{max}
 367 increases when the hurricane intensity ΔP_{max} decreases or the latitude φ increases. Here we
 368 assume that $(\ln R_{max} | \Delta P_{max})$ has the normal distribution proposed by Vickery *et al.* (2000), which
 369 for the region of New Orleans ($\varphi \approx 30^\circ\text{N}$) has parameters

370
$$\begin{aligned} m_{\ln R_{max} | \Delta P_{max}} &= 3.962 - 0.00567 \Delta P_{max} \\ \sigma_{\ln R_{max} | \Delta P_{max}} &= 0.313 \end{aligned} \quad (14)$$

371 where R_{max} is in km and ΔP_{max} is in hPa.

372 The translational speed V_t has weak dependence on the intensity of the TC (Chen *et al.*, 2006;
 373 IPET, 2008) and is usually modeled as a lognormal variable with mean value around 6 m/s and
 374 standard deviation around 2.5 m/s; see Vickery and Twisdale (1995), Vickery *et al.* (2000), and
 375 Chen *et al.* (2006). The former two studies report a slight dependence of V_t on the approach
 376 angle α . To keep the TC recurrence model simple, we use for V_t a lognormal distribution with
 377 the above mean value and standard deviation and assume that V_t and α are independent.

378 Different studies have concluded that the pressure deficit ΔP_{max} has lognormal, Weibull or
 379 Gumbel distribution. The Weibull distribution gives better fits when all tropical cyclones are
 380 considered, whereas the lognormal distribution is more appropriate for storms in the hurricane
 381 intensity range; see Vickery and Twisdale (1995), Chouinard *et al.* (1997) and IPET (2006). The

382 Gumbel distribution has been suggested by IPET (2008) for storms in the CAT35 range ($\Delta P_{max} >$
383 58 hPa). While the Gumbel distribution is appropriate for the analysis of surges, winds and
384 waves (for which the long-term risk is dominated by intense storms), significant rainfall is
385 contributed by less intense slow-moving systems; see Section 5.2 below. For this reason we
386 model ΔP_{max} using the lognormal distribution suggested by IPET (2006). This study shows that
387 for TCs with $\Delta P_{max} > 34$ hPa that made landfall in the longitudinal range 85°-95°W, ΔP_{max} is
388 accurately described by a shifted lognormal distribution with shift parameter 18 hPa, log-mean
389 3.15 and log-standard deviation 0.68.

390 Finally, we set $\lambda = 0.57$ events/year, which is the rate found by IPET (2006) for TCs with
391 $\Delta P_{max} > 34$ hPa making landfall between 85°-95°W along the Gulf of Mexico coast.

392 ***5.2 IDF curves for TC-rainfall and comparison with other storms***

393 Next we use equation (11) with the recurrence model in Section 5.1 to estimate the intensity-
394 duration-frequency (IDF) curves for New Orleans associated with tropical cyclones. The model
395 explicitly accounts for variability in y , V_{max} , R_{max} and V_t . All other input parameters to the MSR
396 model are fixed to the values used in Sections 3 and 4. The joint density of $\{y, V_{max}, R_{max}, V_t\}$ for
397 a TC that makes landfall between longitudes 85°-95°W, $f_{y,\theta}$, is obtained by first calculating the
398 joint density conditional on the pressure deficit ΔP_{max} under the assumption that the variables y ,
399 $[V_{max} | \Delta P_{max}]$, $[R_{max} | \Delta P_{max}]$ and V_t are independent and then averaging the conditional density
400 with respect to ΔP_{max} .

401 Figure 9.a shows the calculated IDF curves as plots of rainfall intensity i against the averaging
402 duration D for different return periods T . For averaging durations below about 12 hours, the
403 decay of i with D follows a power law $D^{-\gamma_D}$ where $\gamma_D \approx 0.55$. This exponent is slightly smaller
404 than the values around 0.6-0.7 that are typical of extra-tropical rainfall (because the rainfall

405 intensities associated with long durations in TCs tend to be high relative to extra-tropical events);
406 see for example Langousis *et al.* (2007). For longer averaging durations, the exponent γ_D rapidly
407 increases and is effectively 1 for $D > 24$ hours; see dashed lines in Figure 9.a. The reason is that
408 the passage of a hurricane usually lasts less than 24 hours; hence for $D > 24$ hours the total
409 rainfall depth is approximately constant and the average rainfall intensity depends on D like D^{-1} .

410 Figure 9.b shows the same results as plots of T against i for different averaging durations D .
411 To determine the importance of TCs relative to other storm types in rainfall risk, the calculated
412 IDF curves are compared with values from TP-40 (Hersfield, 1961), Babak *et al.* (1991) and
413 Singh and Zhang (2007) for return periods $T = 5, 10, 25, 50$ and 100 years. The latter values refer
414 to generic rainfall in the New Orleans area and therefore include both TC and non-TC events.
415 The rainfall values reported in TP-40 cover the whole range of averaging durations D from 0.5-
416 24 hours, whereas Babak *et al.* (1991) and Singh and Zhang (2007) give rainfall values only for
417 $D = 6, 12$ and 24 hours. It is clear from Figure 9.b that for $T > 100$ years also the dependence of
418 the rainfall intensity on T is of the power-law type, say T^{γ_T} with $\gamma_T \approx 0.32$. This exponent is
419 higher than the values around 0.20-0.25 that are typical of ordinary rainfall (Langousis *et al.*,
420 2007; Veneziano *et al.*, 2006b). The higher exponent in tropical cyclones is related to the large
421 dispersion of the amplification factor $\beta_{l,max}$ (see example plots in Figure 8).

422 Another feature of the TC curves in Figure 9.b is the lower asymptote at $T = 1/\lambda = 1.75$ years.
423 This lower bound is a consequence of the fact that the return period of any TC-induced rainfall
424 intensity cannot be lower than the return period of the TCs themselves. The effect of this lower
425 bound is that for short return periods, say $T < 10$ years, the precipitation intensities from tropical
426 cyclones are far below those from ordinary rainfall (frontal events, mesoscale convective
427 systems etc.), for which the recurrence rate is much higher. By contrast, for long averaging

428 durations ($D > 12$ hours) and long return periods ($T = 100$ years), the calculated TC intensities
 429 are close to the empirical intensities, indicating that tropical cyclones have a dominant effect on
 430 those extreme values. Given that the TC curves in Figure 9.b are flatter than those for overall
 431 rain, it is expected that tropical cyclones become even more dominant for longer return periods.

432 For short averaging durations (e.g. D on the order of 1 hour), the contribution of tropical
 433 cyclone rainfall to the risk is negligible, irrespective of the return period. A possible explanation
 434 is that 1) for short averaging durations D , extreme rainfalls are contributed by localized
 435 downpours caused by deep cumulus convection and 2) deep cumulus convection in TCs has
 436 many similarities with tropical cumulus clouds (see e.g. Parrish *et al.*, 1984; Jorgensen *et al.*,
 437 1985; Burpee, 1986; and Powell, 1990 among others). One concludes that for short D rainfall
 438 risk is dominated by storm types whose rate of occurrence is much higher than that of TCs.

439 It is also of interest to determine which tropical cyclones contribute the most to the IDF
 440 values $i(D,T)$. Such TCs might for example be used as scenario events when designing for return
 441 period T . The main parameters to be considered are $\theta = [V_{max}, R_{max}, V_t]$ and the distance y to the
 442 cyclone center. Their modal (most likely) values are obtained by maximizing the conditional
 443 probability density of (y, θ) given $I_{D,max} > i(D,T)$. This conditional density is given by

$$444 \quad f_{y,\theta|D,T}(y, \theta) \propto f_{y,\theta}(y, \theta) P[I_{D,max}(y, \theta) > i(D,T)] \quad (15)$$

445 Figure 10 shows the modal values of V_{max} , R_{max} and V_t for different D and T . The most likely
 446 distance y always satisfies $y \approx R_{max}$. This makes sense because R_{max} is the distance at which the
 447 MSR model predicts maximum large-scale rainfall intensities.

448 Figure 10.a shows that the mode of V_{max} increases when either D or T increase. This makes
 449 physical sense since for any given D , higher rainfall intensities require more intense storms, and
 450 for any given T , intense precipitation over longer averaging durations is associated with more

451 intense systems. Figure 10.b shows that the mode of V_t decreases as T increases, meaning that
452 more intense rainfall is generally produced by slower-moving systems. For averaging durations
453 smaller than 12 hours, the modal value of V_t is insensitive to D , whereas for longer averaging
454 durations V_t decreases faster with T . This faster decay is related to the fact that, for averaging
455 durations D on the order of one day or longer, extremely high rainfall intensities are produced by
456 storms that take a time close to D to pass over the site. Therefore, for T large the translation
457 speed V_t tends to be inversely proportional to D . Finally, Figure 10.c shows that the mode of R_{max}
458 decreases when either D or T increase. This makes sense, since more intense storms tend to have
459 smaller values of R_{max} ; see Section 5.1.

460 **6. Conclusions**

461 We have developed a methodology to assess the frequency of extreme rainfall intensities from
462 tropical cyclones (TCs) in coastal areas with flat topography. The mean rainfall field associated
463 with a TC with maximum tangential wind speed V_{max} , radius of maximum winds R_{max} , and
464 translation speed V_t is obtained using a physically-based (“MSR”) model (Langousis and
465 Veneziano, 2008), whereas rainfall variability at both large scales (from storm to storm) and
466 small scales (due to rainbands and local convection within a single storm) is modeled
467 statistically. The statistical component of the model is estimated using 38 precipitation radar
468 (PR) frames from the TRMM mission; see Table 1. These frames cover a wide range of TC
469 intensities V_{max} and vortex sizes R_{max} . To make the model easier to use in risk analysis, we
470 developed approximate analytical expressions for the statistical parameters. We use Taylor’s
471 hypothesis to convert spatial rainfall intensity fluctuations to temporal fluctuations as the storm
472 passes over a given geographical location A. The combined physical-statistical model predicts
473 the maximum rainfall intensity at A during an averaging period D for a TC with characteristics

474 (V_{max}, R_{max}, V_t) whose center passes at distance y from A. To illustrate the use of the model for
475 long-term rainfall risk analysis, we formulated a recurrence model for tropical cyclones in the
476 Gulf of Mexico that make landfall between longitudes 85° - 95° W and used the rainfall and
477 recurrence models to assess the rainfall risk for New Orleans. Our main findings are as follows.

478 The maximum rainfall $I_{l,max}$ in a spatial interval of length l depends on l , the distance y from
479 the center of the TC, and the intensity V_{max} and size R_{max} of the vortex. We expressed $I_{l,max}$ as the
480 product of the large-scale ($L \approx 400$ km) average rainfall intensity produced by the MSR model,
481 $I_{L,MSR}$, and an amplification factor $\beta_{l,max}$ that includes both storm-to-storm variability and spatial
482 fluctuations of rainfall intensity within a storm. The distribution of $\beta_{l,max}$ depends of course on l ,
483 but in addition depends significantly on the large-scale intensity $I_{L,MSR}$ and the standardized
484 distance from the storm center, $y' = |y/R_{max}|$. Specifically, the dispersion of $\beta_{l,max}$ increases as l
485 and $I_{L,MSR}$ decrease or $y' = |y/R_{max}|$ increases. These trends with $I_{L,MSR}$ and y' are linked to the fact
486 that lower intensity storms and larger distances y' are associated with higher dry area fractions,
487 more intermittent rainfall, and therefore an increased dispersion of the rainfall maxima.

488 Application of the model to TC rainfall risk for New Orleans has produced interesting insight
489 into the importance of tropical cyclones relative to other rainfall-producing events. For short
490 return periods T , the TC intensities are significantly below those from other storms, which have a
491 much higher rate of occurrence. However, as the return period T increases, the TC estimates for
492 long averaging durations (D around 12-24 hours) approach the values found from continuous
493 rainfall records. This means that for long return periods, the long-duration TC rainfalls tend to
494 dominate. In New Orleans, this happens for T around 100 years.

495 To determine how the most likely TC scenario varies with the averaging duration D and the
496 return period T , we calculated the joint distribution of $\{V_{max}, R_{max}, V_t, y\}$ conditioned on

497 exceeding the T -year rainfall intensity for averaging duration D . Then we plotted the modal
498 values of V_{max} , R_{max} , and V_t against D and T ; see Figure 10 (for y , the modal value is always close
499 to R_{max}). The modal value of V_{max} increases when D or T increase, whereas the opposite is true
500 for R_{max} . The mode of the translation velocity V_t is insensitive to D for $D < 24$ hours, but
501 decreases with increasing T and with increasing D for $D > 24$ hours.

502 A rich parameterization and high computational efficiency make the proposed model
503 attractive for rainfall risk applications in TC-prone areas. A limitation of the current model is that
504 it does not account for landfall effects and therefore is applicable only to open-water or coastal
505 sites with flat topography. Future work should focus on extending the model to include inland
506 conditions and extra-tropical conversion using coastal and over-land weather radar data.

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644 Wolff, D.B., D.A. Marks, E. Amitai, D.S. Silberstein, B.L. Fisher, A. Tokay, J. Wang and J.L.
645 Pippitt, 2005: Ground Validation for the Tropical Rainfall Measuring Mission (TRMM), *J.*
646 *Atmos. Oceanic Technol.*, **22**, 365–380.

647 Table 1: Characteristics of the PR/TRMM rain frames used in the analysis. The direction of
648 storm translation is relative to the East and is positive counter-clockwise. The estimates of V_{max}
649 and R_{max} are from the extended best track record (M. DeMaria, 2008; personal communication).

	Storm center		Storm speed (m/s)	storm direction (deg)	V_{max} (m/s)	R_{max} (km)	TRMM frame number	Storm intensity
	Lat. (deg)	Lon. (deg)						
Floyd '99	21.7	-61.6	4.9	143	48.8	41	10290	CAT2
	23.5	-68.7	4.8	169	64.0	37	10317	CAT4
	23.7	-70.6	5.8	171	69.3	37	10321	CAT4
Frances '04	12.6	-43.7	10.9	158	23.1	37	38646	TS
	15.7	-49.8	5.4	139	51.4	19	38667	CAT3
	17	-51.3	5.3	139	54.0	28	38677	CAT3
	17.9	-52.6	4.3	144	59.1	28	38682	CAT4
	19	-57.3	4.9	180	51.4	28	38708	CAT3
	21.2	-68.5	6.1	162	61.7	28	38739	CAT4
Ivan '04	8.9	-38.9	7.6	184	25.7	37	38789	TS
	10.7	-50.6	12.2	185	57.5	28	38814	CAT4
	11.2	-53.4	8.1	173	51.4	28	38820	CAT3
	12.3	-64.1	8.3	166	61.7	19	38845	CAT4
	12.7	-66.2	7.3	164	61.7	20	38851	CAT4
	17.4	-77.3	4.1	194	66.8	28	38892	CAT4
	17.7	-78.4	4.4	153	64.3	28	38897	CAT4
	25.6	-87.4	5.5	112	61.7	46	38954	CAT4
Jeanne '04	27.4	-70.6	5.5	0	38.6	42	39045	CAT1
	25.5	-69.5	1.1	207	41.1	37	39079	CAT2
	26.5	-74.3	7.4	173	43.7	60	39106	CAT2
	26.5	-75.6	6.5	180	46.3	46	39110	CAT2
Karl '04	11.5	-35.3	7.1	176	26.7	37	38987	TS
	17.3	-45.5	2.0	166	57.8	32	39033	CAT3
	19.1	-47.4	5.9	121	64.0	32	39048	CAT4
	22.9	-48.6	8.2	112	54.0	28	39059	CAT3
	25.7	-49.5	6.8	117	48.8	28	39063	CAT3
Katrina '05	24.6	-85.6	2.1	153	51.5	56	44357	CAT3
	25	-86.2	3.5	146	56.5	50	44361	CAT3
	26.9	-89	5.5	135	75.0	38	44373	CAT5
Lilli '02	23.6	-87.2	9.0	162	51.5	20	27826	CAT2
	24.4	-88.4	6.2	141	56.5	20	27830	CAT2
	28.4	-91.4	10.1	117	54.0	20	27842	CAT4
	29	-91.9	5.4	124	41.1	20	27845	CAT2
Rita '05	24.3	-85.9	5.7	189	61.7	28	44743	CAT4
	24.9	-88	3.9	166	77.1	19	44754	CAT5
	25.4	-88.7	4.3	153	72.0	19	44758	CAT5
	26.8	-91	5.5	135	59.1	37	44770	CAT4
	27.4	-91.9	4.8	143	59.1	37	44773	CAT4

650

651

Figure captions

652 Figure 1: Schematic representation of a moving storm. Point O translates with the storm at speed

653 V_t . Point A is the geographical location of interest.

654 Figure 2: Rainfall intensities from Hurricane Katrina (Aug. 28, 2005, at 03:00UTC; TRMM

655 frame 44361) along a cross-section C at distance $y = 100$ km from the storm center, for

656 spatial averaging scales $l = 6$ and 24 km. The maximum values $I_{l,max}$ are indicated by

657 circles. I_L is the average value for the entire cross-section and $I_{L,MSR}$ is the estimate of

658 I_L produced by the MSR model.

659 Figure 3: (a,b) Mean value and standard deviation of $\ln\beta_L$ as a function of the model rainfall

660 intensity $I_{L,MSR}$ and the standardized distance $y' = |y/R_{max}|$ from the TC center using 789

661 cross-sections of the 38 frames in Table 1. The contour plots are obtained using a

662 smoothing Gaussian kernel with standard deviation 0.5. The dashed lines delimit the

663 region of high data density along the direction of the gradient of $\sigma_{\ln\beta_L}$ (white arrow). (c)

664 Plots of $m_{\ln\beta_L}$ and $\sigma_{\ln\beta_L}$ as a function of $\omega = \ln(y') - 0.4\ln(I_{L,MSR})$ along cross-section A.

665 (d) Comparison between the standard normal density and the empirical PDF of $\ln(\beta_L)$,

666 standardized to have zero mean and unit variance.

667 Figure 4: Log-log plots of $E[\gamma_{l,max}]$ and $\text{Var}[\gamma_{l,max}]$ against l for different ranges of I_L . Triangles

668 and circles indicate empirical values. The solid lines are from equation (7).

669 Figure 5: Dependence of the parameters a_1 - a_5 in equation (7) on I_L . The solid lines are least

670 squares fits.

671 Figure 6: Comparison of histograms of $\gamma_{l,max}$ for $l = 96$ and 6 km and different large-scale

672 intensities with theoretical distributions from equations (7) and (8). The intensity

673 categories are the same as in the left column of Figure 4.

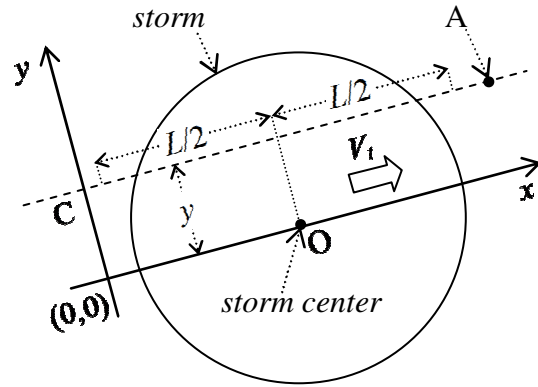
674 Figure 7: Histogram of the non-exceedance probability P in equation (10) for different spatial
675 scales l . Each histogram is based on a sample of size 789.

676 Figure 8: Comparison of the probability density functions of $\beta_{l,max} = I_{l,max}/I_{L,MSR}$ for different V_{max} ,
677 $y' = |y/R_{max}|$, and l .

678 Figure 9: Theoretical IDF curves for New Orleans obtained from equation (11). (a) Maximum
679 rainfall intensity i as a function of averaging duration D for different return periods T .
680 (b) Comparison of the IDF values in (a) for different averaging durations D (solid
681 lines) with intensities obtained from continuous rainfall records.

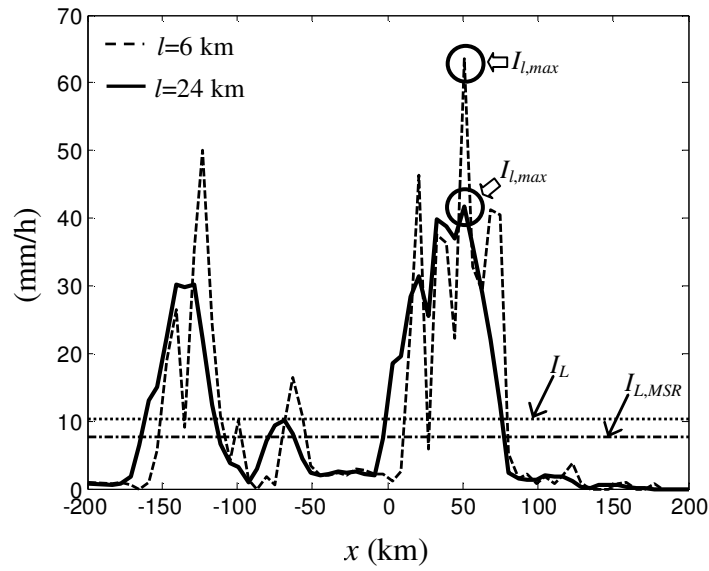
682 Figure 10: Modal values of (V_{max}, V_t, R_{max}) conditioned on exceeding the T -year rainfall intensity
683 for averaging duration $D = 0.5, 1, 3, 6, 12$ and 24 hours.

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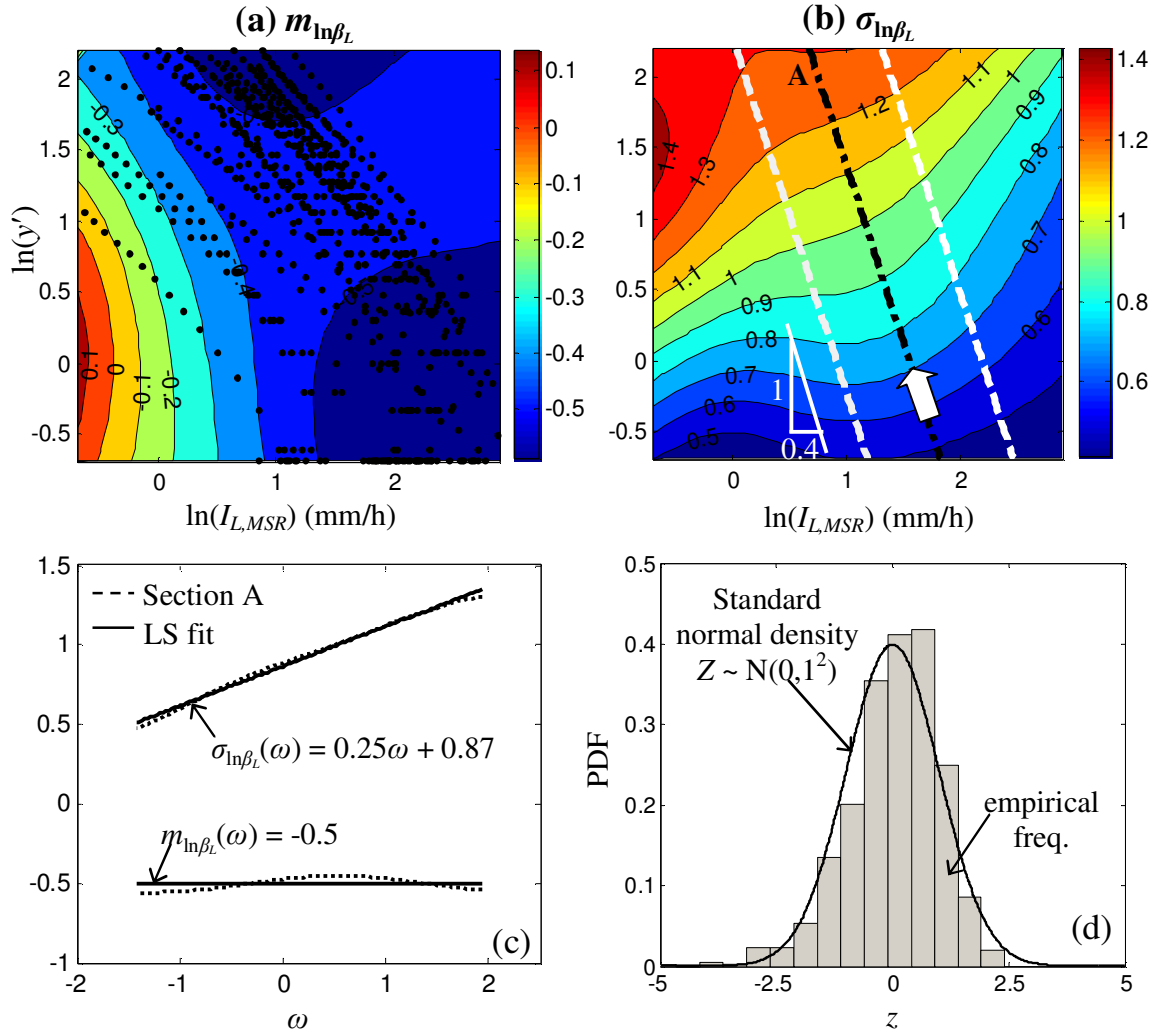
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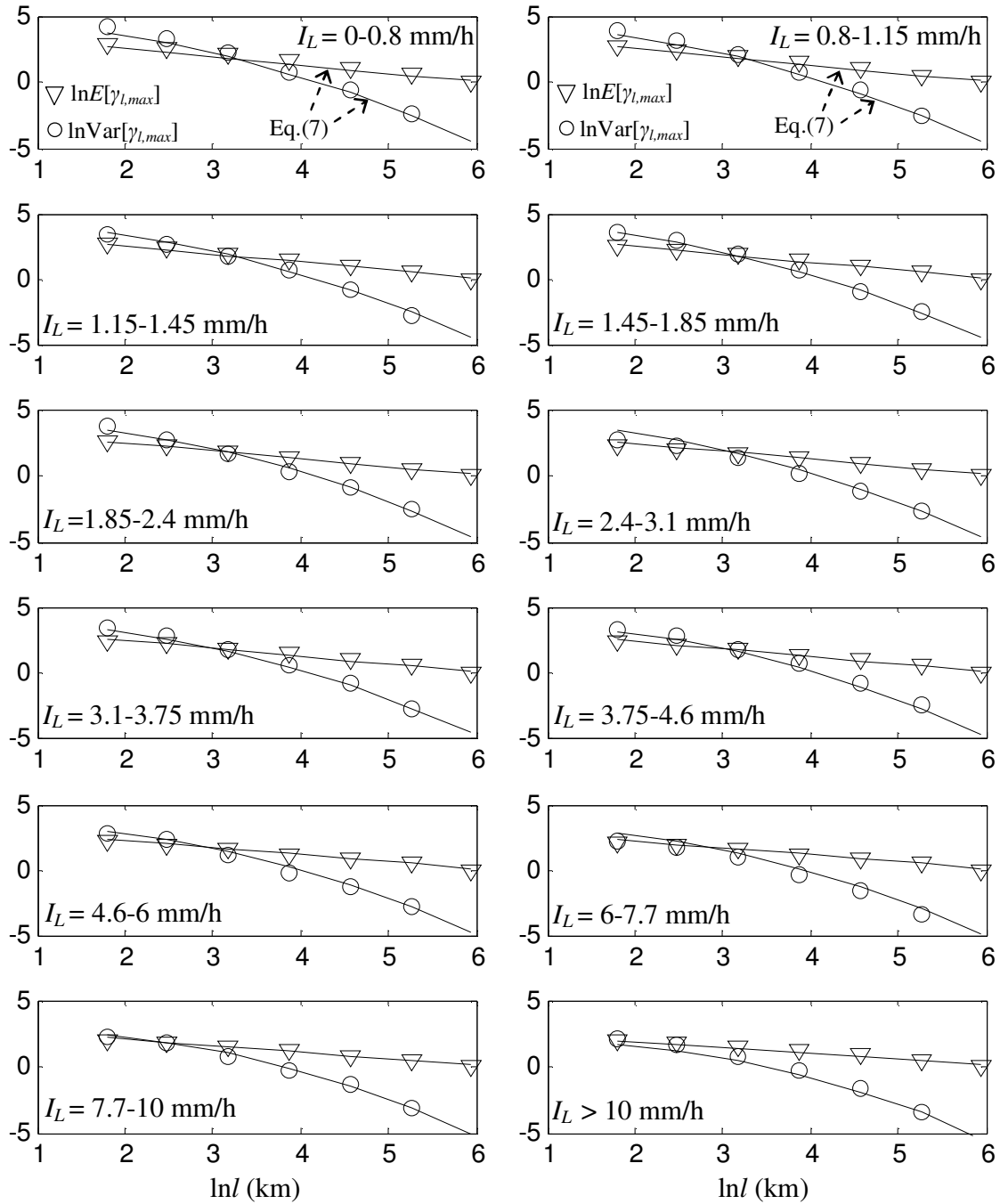
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 692 average value for the entire cross-section and $I_{L,MSR}$ is the estimate of I_L produced by the MSR
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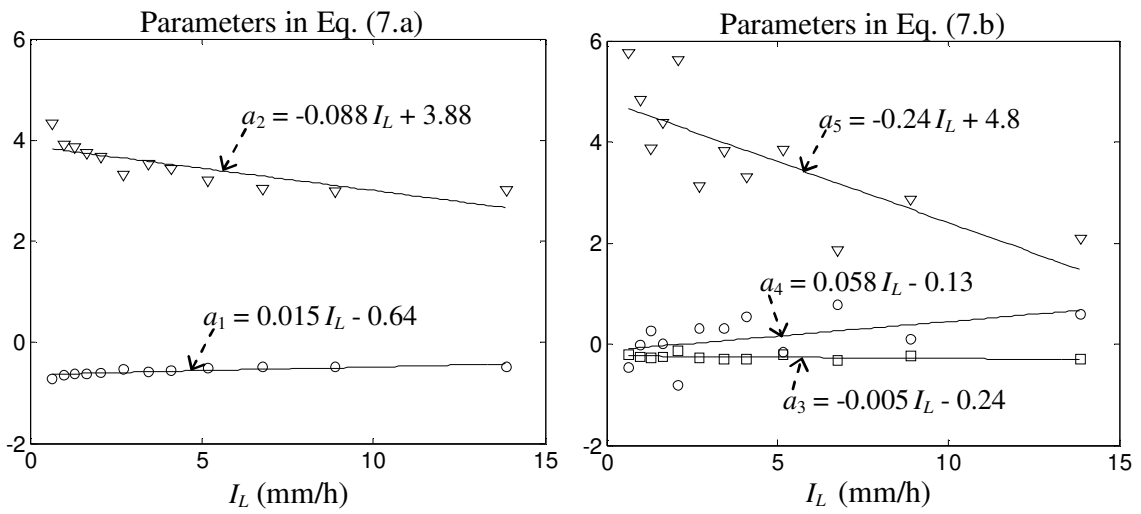
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 702 empirical PDF of $\ln(\beta_L)$, standardized to have zero mean and unit variance.



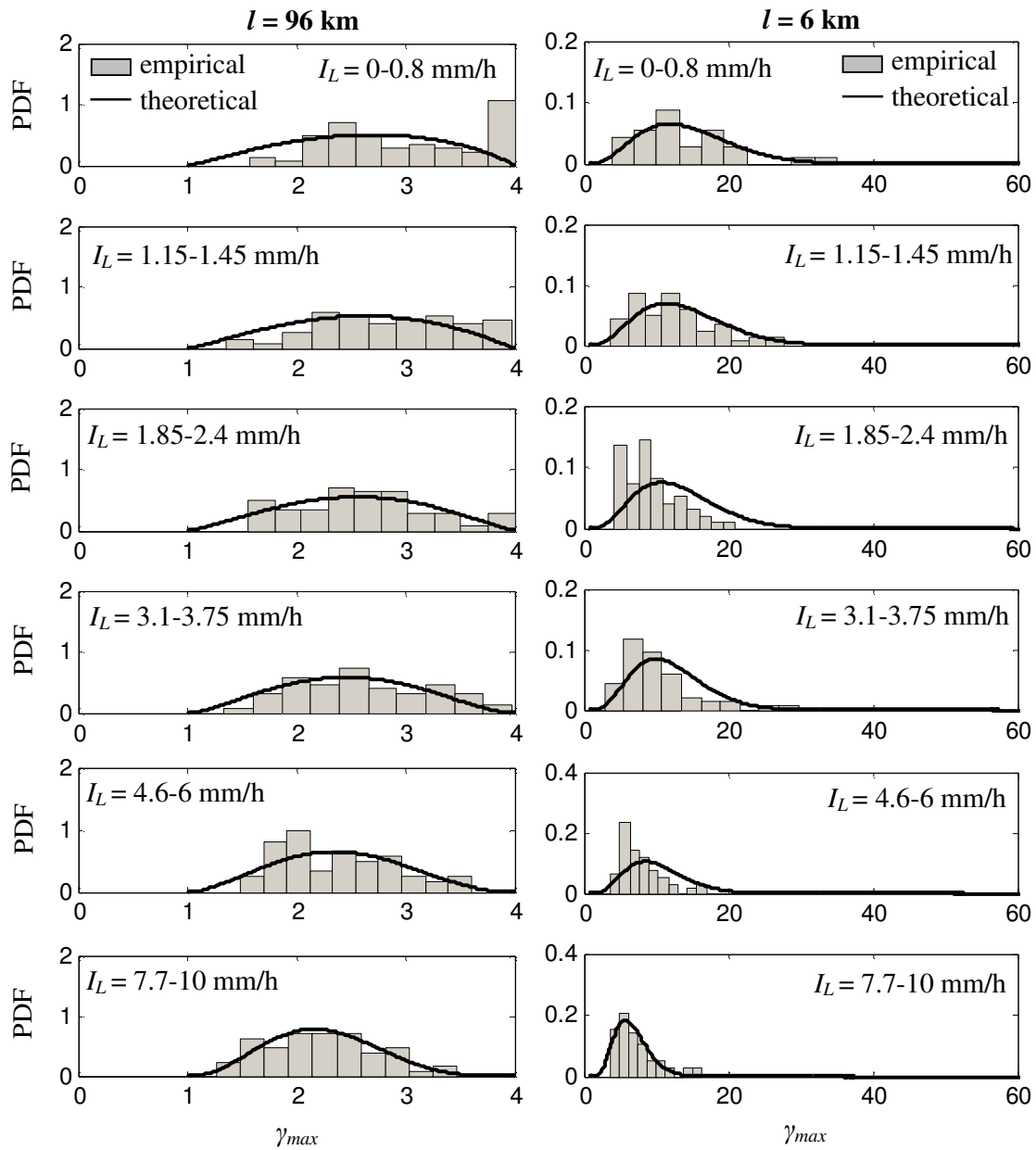
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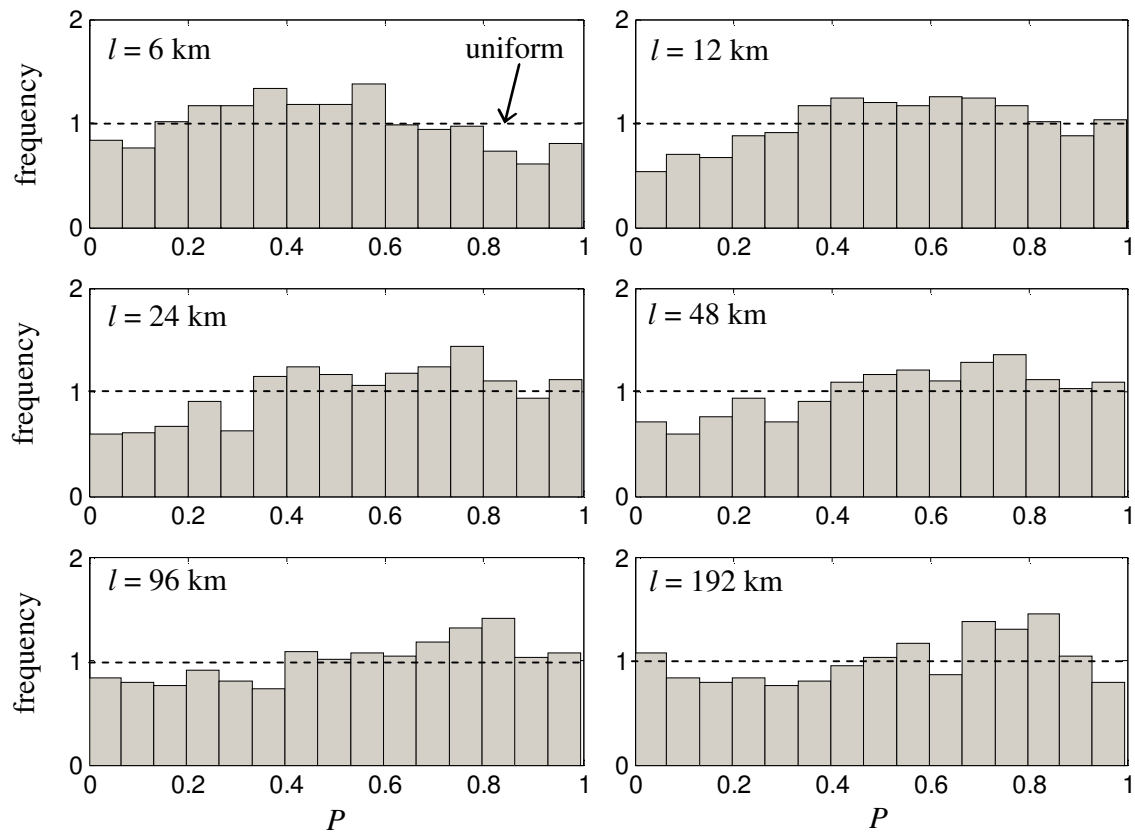
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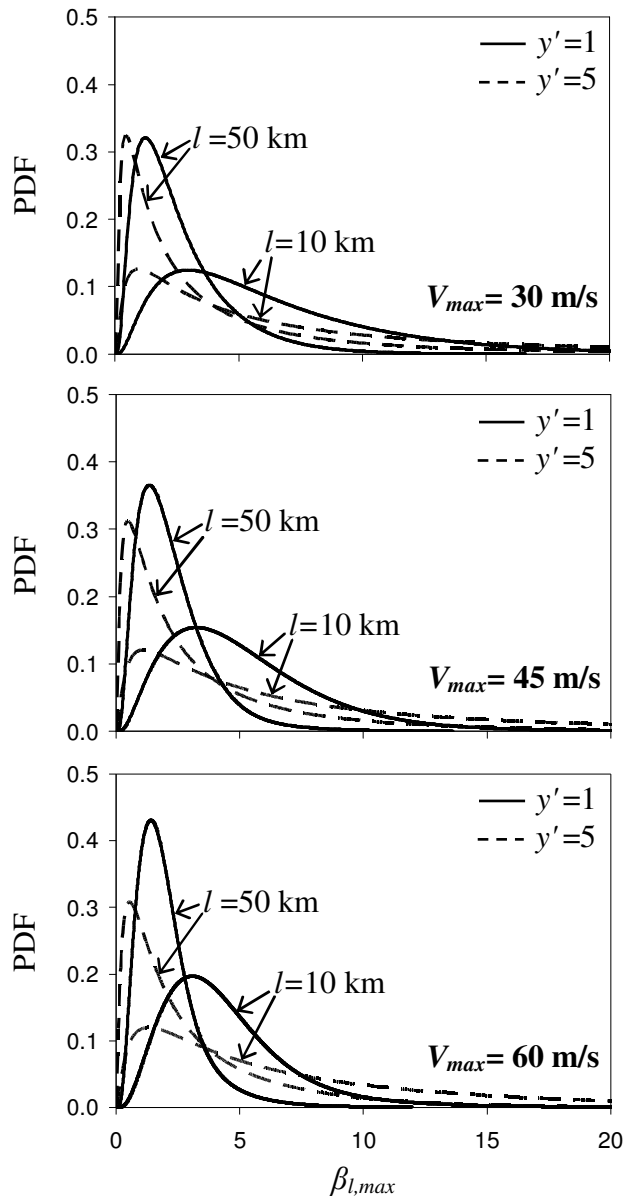
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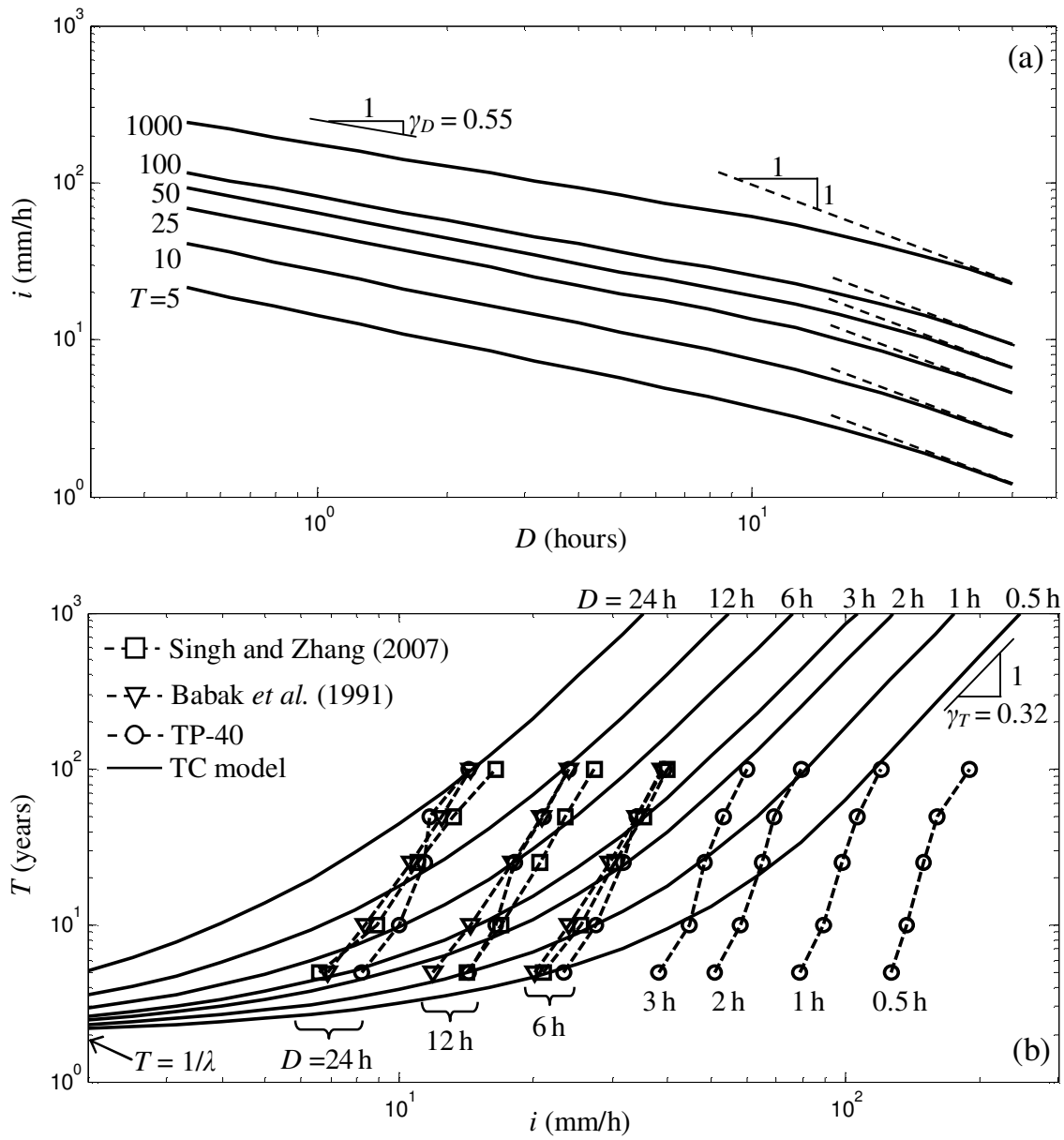
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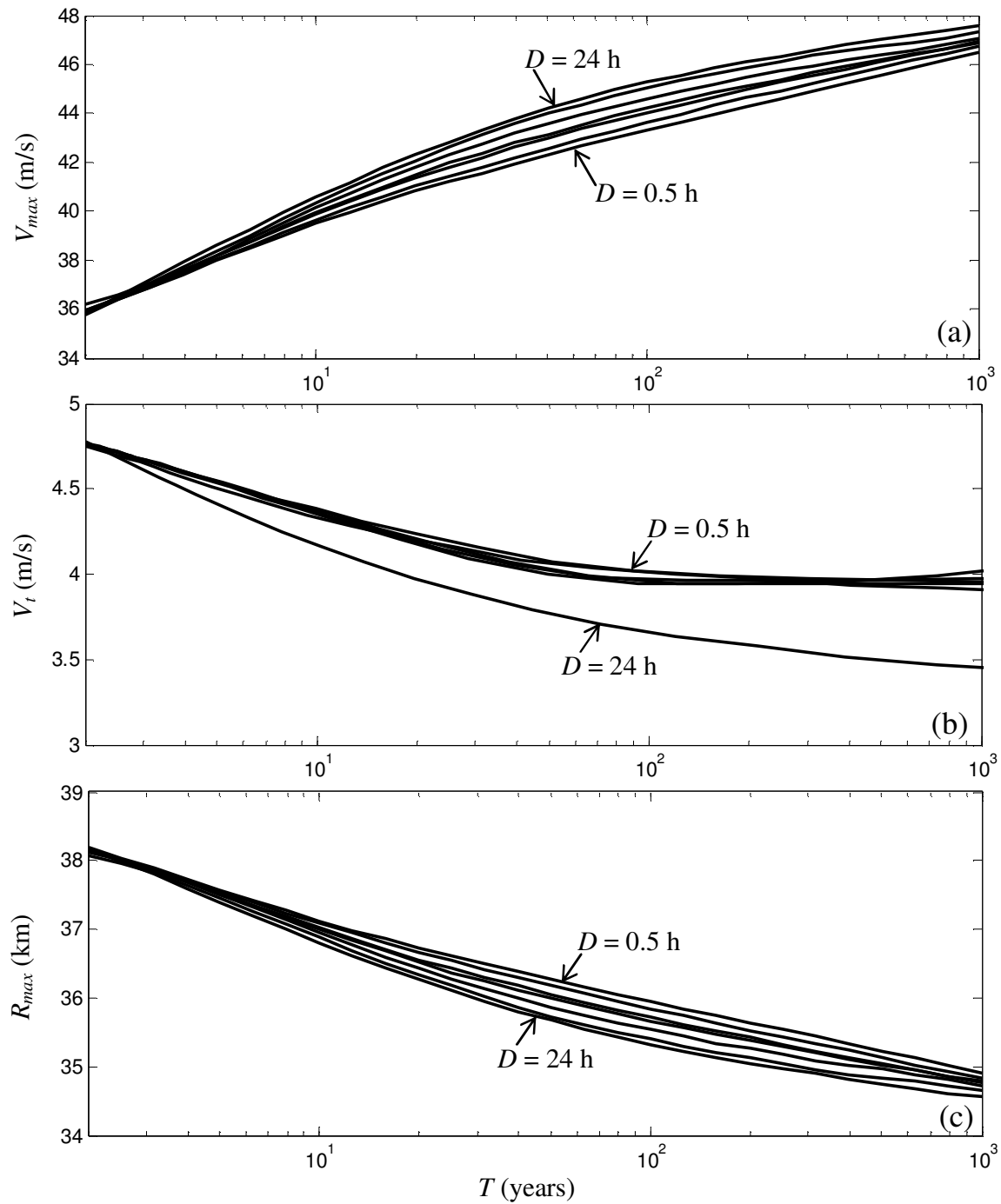
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