### Objective

Find the exact distribution of the maximum of a multifractal cascade and develop simple and accurate approximations.

## 1. Bare cascade construction and dressing

- $\succ$  Start at level 0 with a single *d*-dimensional cubic tile *S* and a unit measure with uniform bare density in S.
- > At each subsequent level n = 1, 2, ..., each tile at the previous level *n* - 1 is partitioned into  $m = \mu^d$  cubic tiles  $S_i$  $(i=1,\ldots,m)$  where  $\mu$  is an integer larger than 1; see Figure 1. The bare density in  $S_i$  is obtained as the product of the bare density in the parent tile at level n-1 and an independent copy  $Y_i$  of a non-negative, unit-mean random variable  $Y_i$ called the cascade generator.





We call *m* the *volumetric multiplicity* of the cascade and r = $m^n$  the volumetric resolution when the cascade has reached level n.

### Dressing

- $\varepsilon_{b,n}$  = bare measure density in a cascade tile at level *n*.  $\varepsilon_{d,n}$  = dressed measure density in a cascade tile at level *n*.
- Z = dressing factor (Z has the same distribution as  $\varepsilon_{d,0}$ ).
- $q^*$  = order above which the moments  $E[Z^q]$  and  $E[(\varepsilon_{d,n})^q]$ diverge.

### The bare and dressed densities satisfy:



### 2. Exact distribution of cascade maxima

Let  $M_n = \max(\varepsilon_{d,n})$  be the maximum in S of the dressed measure density at resolution  $r = m^n$ . Its exact distribution  $F_M$ can be found through the following recursive procedure.  $F_{1} = F_{1}$ n = 0

$$F_{\log M_n} = (F_{\log M_{n-1}}^* f_{\log Y})^m, n = 1, 2, ...$$
(2)

Where  $F_X$  and  $f_X$  denote the CDF and PDF of a random variable X. Notice that calculation of both  $F_{logZ}$  in the first step and  $F_{log_M}$  at each subsequent step requires numerical convolution.

#### Hence need for approximations.



different T.

 $M_{r,T \mid m=4}$ 

 $M_{r,T|m=2}$ 

 $C_{ln}=0.1.$ 

0.80

0.75

3. Effect of the multiplicity *m* 

multiplicity m? If it does not, then:

continuous processes with the same K(q);

 $\succ$  The dimensionality d of the support does not matter.

the maximum at resolution r,  $M_r$ , depend on the volumetric

> Discrete cascade extremes approximate well the extremes of

To answer this question, let  $M_{r,T,m}$  be the upper (1/T)-quantile of

 $M_r$  for a cascade with multiplicity m. For a specific lognormal

cascade, Figure 2 shows the ratio  $M_{r,T,m=4}/M_{r,T,m=2}$  against r for

T: return period expressed in cascade realizations

log<sub>10</sub> /

Figure 2: Ratio  $M_{r,T,m=4}/M_{r,T,m=2}$  for a log normal cascade with

From Figure 2 and similar results for other beta-lognormal

· the correlation of the measure density between tiles at given

• the distribution of the dressing factor Z tightens around 1,

Net effect:  $M_{r,T|m}$  decreases with increasing m.

We explore various approximations to  $M_{\nu}$  in which the

distribution of Z is simplified and the dependence among the

measure densities in different cascade tiles is ignored or treated

in approximation. An important consequence of these

approximations is that convolution operations are avoided.

Below we consider, in sequence, the effect of the following

causing  $M_{r,T|m}$  to decrease. This is a larger effect, especially

distance decreases, producing higher values of  $M_{r,T|m}$ .

cascades one concludes that, as *m* increases:

However this is a small effect.

for small *r* and large *T*.

4. Approximations to  $M_n$ 

simplifications:

Cin: lognormal co-dimension coefficient

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# European Geosciences Union, General Assembly 2005, Vienna Austria, 24-29 April 2005 The maximum of Multifractal Cascades: Exact Distribution and Approximations

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- Assume tile independence. As a part of this analysis, we study the separate effect of short-range and long-range dependences. For a given bare moment scaling function  $K(q) = \log_m[Y^q]$ , does
  - > Approximate Z through a random variable with the distribution type of the bare density  $\varepsilon_{b}$ .
  - $\succ$  Approximate Z and include tile dependence through a simplified correction factor on the quantiles of  $M_{\rm p}$ .

### 4.1 Effect of tile dependence

For a lognormal cascade with m = 2 and co-dimension parameter  $C_{ln} = 0.1$ . Figure 3 compares the distributions of  $M_{n=10}$ and  $M_{n=20}$  with the same distributions under independence. The distribution under independence is found from:





Figure 3: The effect of dependence among cascade tiles on the distribution of  $M_n$ .

- Dependence among the cascade tiles affects the body and lower tail of the distribution of  $M_n$  but not the upper tail.
- > Having found that dependence among the cascade tiles has a significant effect on  $M_n$ , we investigate whether such effect comes mainly from short-range or long-range dependence. **Figure 4** compares the distribution of  $M_{n=20}$  for the cascade in Figure 3 when long-range or short-range dependence is progressively neglected. We conclude that most of the effect comes from long-range dependence.





Figure 4: Change in the distribution of  $M_{n=20}$  when long-range or shortrange dependence is progressively excluded.

### 4.2 Approximation of Z

While possible, numerical calculation of the distribution of Z is tedious (Veneziano and Furcolo, 2003). Here we approximate Z*m* - 1 equals

$$r_{Z} = \left(\frac{m-1}{m-m^{K(2)}}\right)^{1/K(2)}$$
(4)

with the distribution obtained with  $\varepsilon_{b, r_z}$  in place of Z. The distributions are very similar exept for the combination of very low n and and very high quantiles.



Figure 5: Lognormal cascade with  $C_{ln}=0.1$ . Comparison of the exact distribution of  $M_n$  with the distribution obtained by replacing Z with  $\varepsilon_h$ .

### 4.3 Proposed approximation for $M_n$

 $\varkappa$  can be accurately approximated by  $\varepsilon_{b,r,r}$  but tile dependence should not be ignored.

> To approximate the effect of dependence we express the upper (1/T)-quantile of  $M_n, M_{n,T}$ , as:



where  $\gamma(n,T)$  is a numerically evaluated correction factor. For the case of beta-log normal cascades we have found that y(n,T) depends minimally on the beta component.



Figure 6:  $\gamma(n, T)$  as a function of T and n. Lognormal and betalognormal cascades with parameters  $C_{ln} = 0.1$  and  $C_{\beta} = 0, 0.2$ .

We have also found that  $\gamma(n,T)$  depends analytically on the lognormal co-dimension parameter  $C_{ln}$  as:

γc

$$_{ln_1}(n, T) \approx \left(\gamma_{C_{ln_2}}(n, T)\right)^{\sqrt{C_{ln_1}/C_{ln_2}}}$$

This is a special case of the dependence of v(n, T) on the C parameter of log-stable cascades; see Veneziano and Langousis (2005). Insensitivity to the beta component and the above dependence on Cln make Figure 6 sufficient to determine  $\gamma$  for all beta-lognormal cascades.

# 5. Conclusions

- > While possible, numerical calculation of the exact distribution of the cascade maxima  $M_n$  is tedious;
- The cascade multiplicity m has a mild effect on the distribution of  $M_n$  at given volumetric resolution  $r = m^n$ . Therefore, the distribution of the maximum at resolution rdepends little on the dimension *d* of the support;
- > Dependence among the cascade tiles has an important effect on the body and lower tail of the distribution of  $M_n$ ;
- >For beta-log normal cascades, a simple and accurate approximation to  $M_n$  is obtained by combining Eqs. (3), (4) and (5) as follows:

1. Calculate the parameter  $r_Z$  in the approximation of Z;

- 2. Calculate the distribution of  $M_{n ind}$ , the maximum for independent cascade tiles;
- 3. Multiply the upper (1/T)-quantile of  $M_{n,ind}$  by  $\gamma(n,T)$ .

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### References

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with the bare density  $\varepsilon_{b, r_{z}}$  where the resolution  $r_{Z}$  is chosen to match some characteristic of Z. For example, if Z has finite second moment, then one may choose  $r_Z$  so that  $E[(\varepsilon_{b,r_z})^2] = r_Z^{K(2)}$ 

$$E[Z^{2}] = \frac{1}{m - m^{K(2)}}, \text{ giving:}$$

$$r_{Z} = \left(\frac{m - 1}{m - m^{K(2)}}\right)^{1/K(2)}$$

Figure 5 compares the upper tail of the exact distribution of  $M_n$